



Зондовая микроскопия: методы, теория, приложения

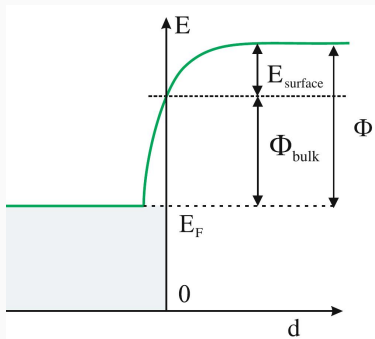
Лекция 9: Кельвиновская микроскопия
Сканирующая туннельная микроскопия

О.В. Синицына, Г.Б. Мешков, Я.В. Гиндикин

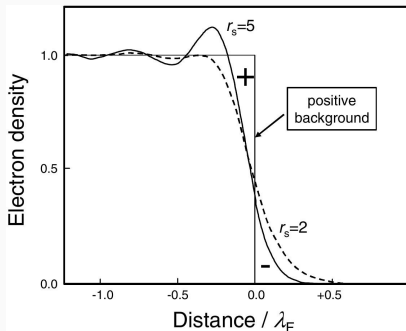
16 апреля 2018г

Московский государственный университет
Факультет наук о материалах

Зонная диаграмма

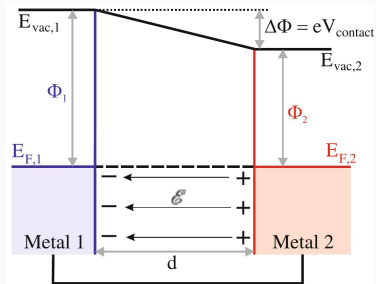
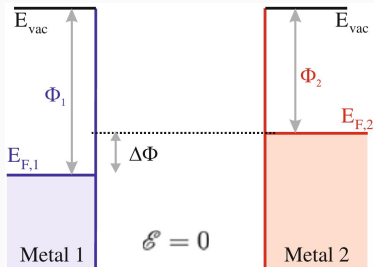


Эффект Смолуховского

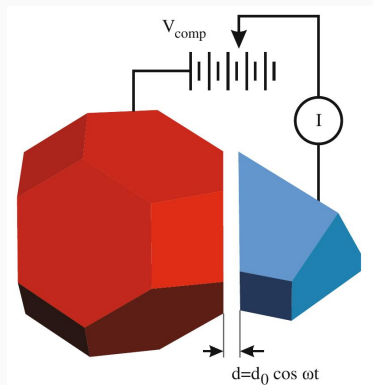


$$\Phi(d) = E_{\text{out}}(d) - E_F$$

Контактная разность потенциалов



Метод Кельвина



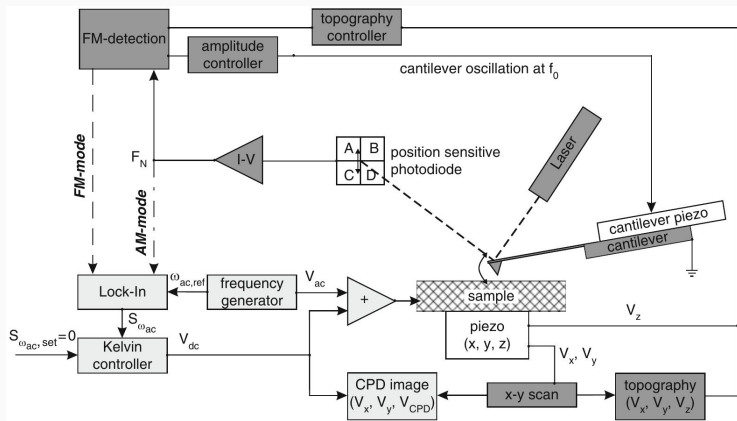
$$Q = CV = C(V_{\text{contact}} - V_{\text{comp}})$$

$$I = \frac{dQ}{dt} = \frac{dC}{dt}(V_{\text{contact}} - V_{\text{comp}})$$

$$I \rightarrow 0$$

$$V_{\text{comp}} = V_{\text{contact}} = \frac{1}{e} \Delta \Phi$$

Кельвиновская микроскопия

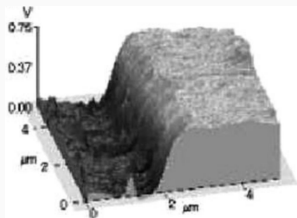
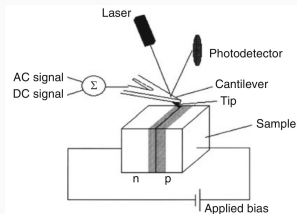


$$F = -\frac{1}{2} \frac{\partial C}{\partial z} (V_{dc} - V_{CPD} + V_{ac} \sin(\omega_{ac} t))^2; \quad F_{2\omega_{ac}} = \frac{\partial C}{\partial z} \frac{V_{ac}^2}{4} \cos(2\omega_{ac} t)$$

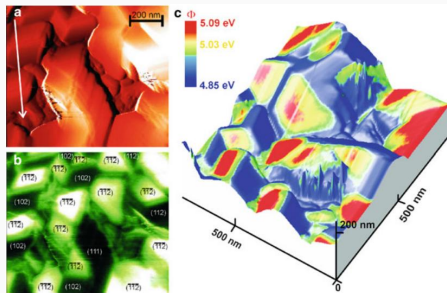
$$F_{\omega_{ac}} = -\frac{\partial C}{\partial z} (V_{dc} - V_{CPD}) V_{ac} \sin(\omega_{ac} t); \quad F_{dc} = -\frac{\partial C}{\partial z} \left(\frac{1}{2} (V_{dc} - V_{CPD})^2 + \frac{V_{ac}^2}{4} \right)$$

Кельвиновская микроскопия: примеры

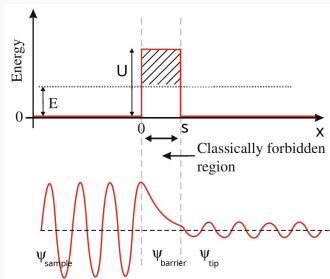
GaP p-n переход



Пленка CuGaSe_2 на подложке $\text{ZnSe}(110)$



Туннелирование: наивный подход



$$-\frac{\hbar^2}{2m}\partial_x^2\psi(x) + U(x)\psi(x) = E\psi(x)$$

$$\psi_{\text{sample}} = e^{ikx} + Ae^{-ikx}, \quad k = \frac{\sqrt{2mE}}{\hbar}$$

$$\psi_{\text{barrier}}(x) = Be^{-\kappa x} + Ce^{\kappa x},$$

$$\kappa = \frac{\sqrt{2m(U-E)}}{\hbar}$$

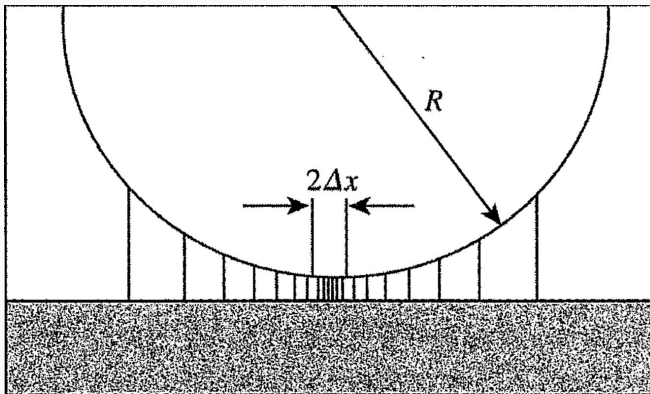
$$\psi_{\text{tip}} = De^{ikx}$$

$$I = -\frac{i\hbar}{2m}(\psi^*\partial_x\psi - \psi\partial_x\psi^*)$$

$$T = \frac{I_{\text{tip}}}{I_{\text{sample}}} = |D|^2 \sim \frac{16\kappa^2 k^2}{(\kappa^2 + k^2)^2} e^{-2\kappa s}, \quad \kappa s \gg 1$$

$$\kappa = \frac{\sqrt{2m\Phi}}{\hbar} \approx 1\text{\AA}^{-1}, \quad \Phi \approx 5\text{eV}$$

СТМ: оценка разрешения

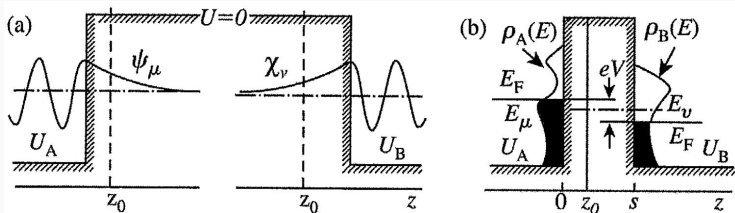


$$\Delta z \approx \frac{\Delta x^2}{2R}$$

$$I(\Delta x) \approx e^{-\kappa \frac{\Delta x^2}{R}}$$

$$R = 10\text{\AA}, \quad \Delta x \approx 3\text{\AA}$$

СТМ: теория Бардина



$$i\hbar\partial_t\Psi = \left[-\frac{\hbar^2}{2m}\partial_z^2 + U_A\right]\Psi,$$

$$\Psi = \psi_\mu e^{-iE_\mu t/\hbar},$$

$$\left[-\frac{\hbar^2}{2m}\partial_z^2 + U_A\right]\psi_\mu = E_\mu\psi_\mu,$$

$$i\hbar\partial_t\Psi = \left[-\frac{\hbar^2}{2m}\partial_z^2 + U_B\right]\Psi$$

$$\Psi = \chi_\nu e^{-iE_\nu t/\hbar}$$

$$\left[-\frac{\hbar^2}{2m}\partial_z^2 + U_B\right]\chi_\nu = E_\nu\chi_\nu$$

$$i\hbar\partial_t\Psi = \left[-\frac{\hbar^2}{2m}\partial_z^2 + U_A + U_B \right] \Psi$$

$$\Psi = \psi_\mu e^{-iE_\mu t/\hbar} + \sum_{\nu=1}^{\infty} c_\nu(t) \chi_\nu e^{-iE_\nu t/\hbar}, \quad c_\nu(0) = 0$$

$$\int \psi_\mu^* \chi_\nu d^3r \approx 0 \leftarrow \text{модельное предположение}$$

$$i\hbar\partial_t c_\nu = \int_{z>z_0} \psi_\mu U_B \chi_\nu^* d^3r \cdot e^{-i(E_\mu - E_\nu)t/\hbar}$$

$$M_{\mu\nu} = \int_{z>z_0} \psi_\mu U_B \chi_\nu^* d^3r \leftarrow \text{туннельный матричный элемент}$$

$$c_\nu(t) = M_{\mu\nu} \frac{e^{-i(E_\mu - E_\nu)t/\hbar} - 1}{E_\mu - E_\nu}$$

$$p_{\mu\nu}(t) = |c_\nu(t)|^2 = |M_{\mu\nu}|^2 \frac{4 \sin^2[(E_\mu - E_\nu)t/2\hbar]}{(E_\mu - E_\nu)^2}$$

$$\lim_{t \rightarrow +\infty} \frac{1}{\omega^2 t} \sin^2(\omega t) = \pi \delta(\omega)$$

$$P_{\mu\nu}(t) = \partial_t |c_\nu(t)|^2 = \frac{2\pi}{\hbar} |M_{\mu\nu}|^2 \delta(E_\mu - E_\nu) \leftarrow \text{золотое правило Ферми}$$

$$I_{A \rightarrow B} = \frac{4\pi e}{\hbar} \sum_{\mu\nu} f(E_\mu - E_F^A) (1 - f(E_\nu - E_F^B)) |M_{\mu\nu}|^2 \delta(E_\mu - E_\nu)$$

$$I_{B \rightarrow A} = \frac{4\pi e}{\hbar} \sum_{\mu\nu} f(E_\nu - E_F^B) (1 - f(E_\mu - E_F^A)) |M_{\mu\nu}|^2 \delta(E_\mu - E_\nu)$$

$$I = \frac{4\pi e}{\hbar} \int_{-\infty}^{\infty} (f(E_F - eV + \epsilon) - f(E_F + \epsilon)) \rho_A(E_F - eV + \epsilon) \rho_B(E_F + \epsilon) |M(\epsilon)|^2 d\epsilon$$

$$I = \frac{4\pi e}{\hbar} \int_0^{eV} \rho_A(E_F - eV + \epsilon) \rho_B(E_F + \epsilon) |M(\epsilon)|^2 d\epsilon \quad \text{при } T \rightarrow 0$$

$$M_{\mu\nu} = \int_{z>z_0} \psi_\mu(E_\nu + \frac{\hbar^2}{2m}\partial_z^2)\chi_\nu^* d^3r$$

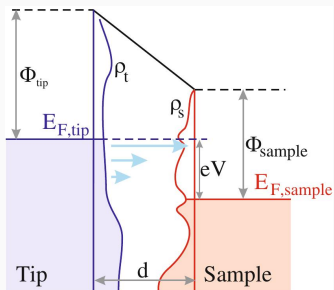
Для случая упругого туннелирования ($E_\mu = E_\nu$),

$$\begin{aligned} M_{\mu\nu} &= \int_{z>z_0} (\chi_\nu^* E_\mu \psi_\mu + \psi_\mu \frac{\hbar^2}{2m} \partial_z^2 \chi_\nu^*) d^3r \\ &= -\frac{\hbar^2}{2m} \int_{z>z_0} (\chi_\nu^* \partial_z^2 \psi_\mu - \psi_\mu \partial_z^2 \chi_\nu^*) d^3r. \end{aligned}$$

$$\chi_\nu^* \partial_z^2 \psi_\mu - \psi_\mu \partial_z^2 \chi_\nu^* = \partial_z (\chi_\nu^* \partial_z \psi_\mu - \psi_\mu \partial_z \chi_\nu^*)$$

$$M_{\mu\nu} = \frac{\hbar^2}{2m} \int_{z=z_0} (\psi_\mu \partial_z \chi_\nu^* - \chi_\nu^* \partial_z \psi_\mu) dx dy$$

СТМ: модель плоского интерфейса



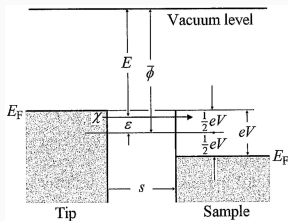
$$\psi_{\mu}(z) = \psi_{\mu}(0)e^{-\kappa_{\mu}z}, \quad \kappa_{\mu} = \sqrt{2m|E_{\mu}|/\hbar}$$

$$\chi_{\nu}(z) = \chi_{\nu}(s)e^{\kappa_{\nu}(z-s)}, \quad \kappa_{\nu} = \kappa_{\mu}$$

$$M_{\mu\nu} = \frac{\hbar^2}{2m} \int_{z=z_0} 2\kappa_{\mu}\psi_{\mu}(0)\chi_{\nu}(s)e^{-\kappa_{\mu}z_0}e^{\kappa_{\mu}(z_0-s)}dx dy$$

$$= \left[\frac{\hbar^2}{2m} \int_{z=z_0} 2\kappa_{\mu}\psi_{\mu}(0)\chi_{\nu}(s)dx dy \right] e^{-\kappa_{\mu}s}.$$

СТМ: асимметрия туннельного спектра



$$I = \frac{4\pi e}{\hbar} \int_{-\frac{1}{2}eV}^{\frac{1}{2}eV} \rho_S(E_F + \frac{1}{2}eV + \epsilon) \rho_T(E_F - \frac{1}{2}eV + \epsilon) |M(\epsilon)|^2 d\epsilon$$

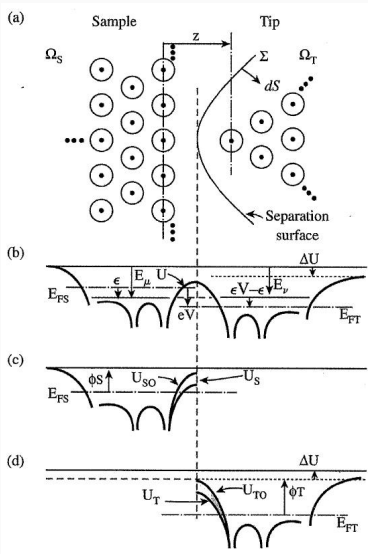
$$\kappa = \frac{\sqrt{2m(\bar{\phi} - \epsilon)}}{\hbar} \approx \kappa_0 \left(1 - \frac{\epsilon}{2\bar{\phi}}\right); \quad \kappa_0 = \frac{\sqrt{2m\bar{\phi}}}{\hbar}$$

$$M(\epsilon) = M(0) \exp\left(\frac{\kappa_0 \epsilon s}{2\bar{\phi}}\right)$$

$$I = \frac{4\pi e}{\hbar} \int_{-\frac{1}{2}eV}^{\frac{1}{2}eV} \rho_S(E_F + \frac{1}{2}eV + \epsilon) \rho_T(E_F - \frac{1}{2}eV + \epsilon) |M(0)|^2 \times \exp\left(\frac{\kappa_0 \epsilon s}{\bar{\phi}}\right) d\epsilon.$$

$$\kappa_0 = 10 \text{ nm}^{-1}, \quad \bar{\phi} = 5 \text{ eV}, \quad s = 1 \text{ nm}, \quad \epsilon = 0.5 \text{ eV}$$

СТМ: модель 3D интерфейса



$$U_S + U_T = U, \quad U_S U_T = 0$$

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + U_S \right] \psi_\mu = E_\mu \psi_\mu$$

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + U_T \right] \chi_\nu = E_\nu \chi_\nu$$

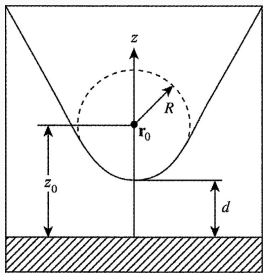
$$i\hbar \partial_t \Psi = \left[-\frac{\hbar^2}{2m} \nabla^2 + U_S + U_T \right] \Psi$$

$$M_{\mu\nu} = \int_{\Omega_T} \psi_\mu U_T \chi_\nu^* d^3r$$

$$M_{\mu\nu} = \frac{\hbar^2}{2m} \int_{\Omega_T} (\chi_\nu^* \nabla^2 \psi_\mu - \psi_\mu \nabla^2 \chi_\nu^*) d^3r$$

$$M_{\mu\nu} = \frac{\hbar^2}{2m} \int_{\Sigma} (\psi_\mu \nabla \chi_\nu^* - \chi_\nu^* \nabla \psi_\mu) \cdot dS = M_{\nu\mu}^*$$

СТМ: модель Терсоффа-Хаманна



$$-\frac{\hbar^2}{2m}\Delta\psi(\mathbf{r}) = -\phi\psi(\mathbf{r})$$

$$\Delta\psi(\mathbf{r}) = \kappa^2\psi(\mathbf{r}), \quad \kappa = \sqrt{2m\phi}/\hbar$$

$$\psi(\mathbf{r}) = \int d^2\mathbf{q} f(\mathbf{q}, z) e^{i\mathbf{q}\cdot\mathbf{x}},$$

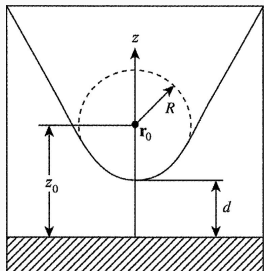
$$\mathbf{q} = (k_x, k_y), \quad \mathbf{x} = (x, y)$$

$$\partial_z^2 f(\mathbf{q}, z) = (\mathbf{q}^2 + \kappa^2) f(\mathbf{q}, z)$$

$$f(\mathbf{q}, z) = a(\mathbf{q}) e^{-\sqrt{\mathbf{q}^2 + \kappa^2} z}$$

$$\psi(\mathbf{r}) = \int d^2\mathbf{q} a(\mathbf{q}) e^{-\sqrt{\mathbf{q}^2 + \kappa^2} z + i\mathbf{q}\cdot\mathbf{x}}$$

СТМ: модель Терсоффа-Хаманна



$$r^2 = x^2 + y^2 + (z - z_0)^2$$

$$\frac{1}{r} \frac{d^2}{dr^2} [r\chi(r)] = \kappa^2 \chi(r)$$

$$\chi(r) = \frac{1}{r} e^{-\kappa r} = \frac{1}{2\pi} \int d^2\mathbf{p} \frac{e^{-\sqrt{\mathbf{p}^2 + \kappa^2}(z_0 - z) + i\mathbf{p} \cdot \mathbf{x}}}{\sqrt{\mathbf{p}^2 + \kappa^2}}$$

$$M = \frac{\hbar^2}{2m} \int_{z=0} (\psi \partial_z \chi^* - \chi^* \partial_z \psi) dx dy$$

$$M \propto \int d^2\mathbf{q} d^2\mathbf{p} d^2\mathbf{x} \left[1 + \frac{\sqrt{\mathbf{q}^2 + \kappa^2}}{\sqrt{\mathbf{p}^2 + \kappa^2}} \right] a(\mathbf{q}) e^{-\sqrt{\mathbf{q}^2 + \kappa^2} z_0 + i(\mathbf{q} + \mathbf{p}) \cdot \mathbf{x}}$$

$$M \propto \int d^2\mathbf{q} a(\mathbf{q}) e^{-\sqrt{\mathbf{q}^2 + \kappa^2} z_0} = \psi(\mathbf{r}_0)$$

$$G \propto |\psi(\mathbf{r}_0)|^2 \rho_S(E_F) = \rho_S(E_F, \mathbf{r}_0)$$



C. J. Chen.

Introduction to scanning tunneling microscopy.

Oxford University Press, 2008.



S. Sadewasser and T. Glatzel.

Kelvin probe force microscopy: measuring and compensating electrostatic forces.

Springer, 2011.



B. Voigtländer.

Scanning probe microscopy: Atomic force microscopy and scanning tunneling microscopy.

Springer, 2015.

Вопросы?